

Higgs bundles and related topics.

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and Abelian Hodge correspondence.
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Lecture 1:

Plan of
the course

Part I: Basics of Higgs bundles
geometry of moduli space

NAH
Higher Teichmüller theory
parabolic Higgs bundles

6

Part II: topics.

6.

Today: Explain Betti, de Rham, Dolbeault
moduli spaces.

Take a close look at rk 1 case.

Reference:

W. Goldman and E.Z. Xia,

"Rank one Higgs bundles and representations of
fundamental groups of R.S."

§0. Equivalence of deformation theories.

Defn. A deformation theory (or transformation groupoid)
 (S, G) consisting of a category \mathcal{C} defined by
a grp action as follows:

Let $\alpha: G \times S \rightarrow S$ left action.

(S, G) consists of the category \mathcal{C} with $\text{Obj}(\mathcal{C}) = S$
with morphism $x \xrightarrow{g} y$ corresponding to

- the triple $(y, x, y) \in G \times S \times S$ s.t. $\alpha(y, x) = y$.
- $e \in G$ determines the identity morphism $x \xrightarrow{e} x$.
- $x \xrightarrow{g} y$ has an inverse $y \xrightarrow{g^{-1}} x$
- composition.

Defn. The moduli set corresponding to such a groupoid is the set $\text{Iso}(\mathcal{L})$ of isomorphism classes of objects.

Defn. An equivalence of categories is a functor $F: \mathcal{A} \rightarrow \mathcal{B}$ s.t. $\exists H: \mathcal{B} \rightarrow \mathcal{A}$ and $F \circ H \cong I_{\mathcal{B}}$
 $H \circ F \cong I_{\mathcal{A}}$.

\rightsquigarrow a bijection: $\text{Isom}(\mathcal{A}) \rightarrow \text{Isom}(\mathcal{B})$.

Prop (Criterion) A functor $F: \mathcal{A} \rightarrow \mathcal{B}$ is an equiv iff

- (1) surjective on Isomorphism classes.
- (2) Full: $F(x, y) = \text{Mor}(x, y) \rightarrow \text{Mor}(F(x), F(y))$ is surjective.
- (3) Faithful: injective.

§1. The Betti groupoid.

Fix G a structure grp, e.g. $GL(n, \mathbb{C}), SL(n, \mathbb{C}), U(n)$
 Σ a compact smooth oriented surface with fundamental grp π .

- The objects are representations: $\pi \rightarrow G$
 $S = \text{Hom}(\pi, G)$
- The morphisms are from G by conjugation.
 $G \times \text{Hom}(\pi, G) \rightarrow \text{Hom}(\pi, G)$
 $g \cdot \rho \mapsto g^{-1} \rho g$

Defn. The Betti groupoid is $(\text{Hom}(\pi, G), G)$.

• π admits a presentation

$$\langle A_1, B_1, \dots, A_g, B_g \mid [A_1, B_1] \dots [A_g, B_g] = 1 \rangle$$

The map $\text{Hom}(\pi, G) \hookrightarrow G^{2g}$

$$p \longmapsto (p(A_1), p(B_1), \dots, p(A_g), p(B_g)).$$

embeds $\text{Hom}(\pi, G)$ as a Zariski-closed subset of G^{2g} defined $[a_1, \beta_1] \dots [a_g, \beta_g] = 1$. (*)

• If G is abelian, it acts trivially on $\text{Hom}(\pi, G)$.

The condition (*) is automatically satisfied.

$$\text{So } \text{Hom}(\pi, G)/G \cong \text{Hom}(\pi, G) \cong G^{2g}.$$

"
Isom($(\text{Hom}(\pi, G), G)$).

will apply this to $G = \mathbb{C}^*, U(1), \mathbb{R}^\pm$.

§2. The de Rham groupoid

Let E be a smooth complex vector bundle over Σ .

$\mathcal{A}^k(\Sigma)$ denote the space of k -forms on Σ

$\mathcal{A}^k(\Sigma, E)$ E -valued k -forms.

Defn. A gauge transformation of E is a smooth bundle automorphism

$$\xi: E \rightarrow E$$

$$\downarrow \Omega \downarrow$$

$$\text{id}: \Sigma \rightarrow \Sigma$$

Denote by $\mathcal{G}(E)$ the group of gauge transformations of E .

Defn. (Connection)

A connection on E is an operator

$$D: \mathcal{A}^0(\Sigma; E) \rightarrow \mathcal{A}^1(\Sigma, E)$$

$$\text{s.t. } D(fs) = fD(s) + df \wedge s.$$

Such a map extends to $D: \mathcal{A}^p(\Sigma; E) \rightarrow \mathcal{A}^{p+1}(\Sigma, E)$.

Denote by $\mathcal{U}(E)$ the space of all connections on E .

Note that fix a connection D_0 , an arbitrary connection

$$D = D_0 + \eta \quad \text{for } \eta \in \mathcal{A}^1(\Sigma; \text{End}(E)).$$

So $\mathcal{U}(E)$ is an affine space modeled on $\mathcal{A}^1(\Sigma; \text{End}(E))$.

Defn. (Curvature)

The curvature of a connection D is

$$\text{defined as } F(D)s = D \circ D(s),$$

turns out to be an $\text{End}(E)$ -valued 2-form

$$F(D) \in \mathcal{A}^2(\Sigma; \text{End}(E)).$$

Call D flat if $F(D) = 0$.

Denote by $\mathcal{F}(E)$ the space of flat connections on E .

(Note that for the existence of a flat connection, require $\text{deg}(E) = 0$.)

• The gauge action on connections

$$\xi^* D \text{ is defined as } (\xi^* D)(s) = \int \xi^* D(\xi \cdot s)$$

for $\xi \in \mathcal{G}(E)$.

$$\xi \cdot D := (\xi^{-1})^* D.$$

Locally, w.r.t a frame e ,

$$D = d + \eta \quad (\text{i.e. } D e = e \eta)$$

$$\text{Then } \xi^* D = d + g^{-1} \eta g + g^{-1} dg$$

(Here, g is the local expression of ξ w.r.t e)

$$\text{i.e. } \xi e = e g.$$

$$\begin{aligned} (\xi^* D)(e) &= D(e g) = e(\eta g + dg) \\ &= e g (g^{-1} \eta g + g^{-1} dg). \end{aligned}$$

$$\bullet F(\xi^* D) = \xi^*(F(D)).$$

Hence, $G(E)$ preserves flatness.

Defn. The de Rham groupoid is $(F(E), G(E))$.

§3. Equivalence between Betti and de Rham groupoids

Start from a flat connection D on a vector bundle E ,
want to obtain a rep $\rho: \pi \rightarrow GL(n, \mathbb{C})$.

Locally, w.r.t a frame e , $D e = e \cdot \eta$.

Over a smooth path $\sigma: [0, 1] \rightarrow \Sigma$,
parallel transport defines a linear map between
the fibers $P_{\sigma(t)}: E_{\sigma(0)} \rightarrow E_{\sigma(t)}$.

That is, $P_{\sigma(t)}(v)$ is parallel w.r.t D , for $v \in E_{\sigma(0)}$.

Suppose $v = (e \circ \gamma(t_0)) \cdot f(t_0) \in E_{\gamma(t_0)}$.

Then $P_{\gamma(t)}(v) = (e \circ \gamma(t)) \cdot \underline{g(t)} \cdot f(t_0)$ is parallel to D

$$\Leftrightarrow D_{\frac{d}{dt}}((e \circ \gamma(t)) \cdot g(t) \cdot f(t_0)) = 0 \quad (\partial/\partial t)$$

$$\Leftrightarrow (e \circ \gamma(t)) \cdot \left(\eta \circ \gamma(t) \cdot g(t) + dg(t) \right) \cdot f(t_0) = 0$$

$$\Leftrightarrow g'(t) + (\eta \circ \gamma(t)) \cdot g(t) = 0$$

$$\Leftrightarrow g(t) = \exp\left(-\int_0^t \gamma^* \eta\right)$$

Fact: Flatness of D implies the parallel transport only depends on homotopic class of γ relative to its endpoints.

Now we obtain a homomorphism: fix a pt $p \in E_{x_0}$.

$$\text{hol}_p(D) : \pi_1(X, x_0) \rightarrow GL(n, \mathbb{C})$$

$$\gamma \longmapsto \left(P_\gamma : E_{x_0} \rightarrow E_{x_0} \right)^{-1}$$

w.r.t a fixed frame e at E_{x_0} .

Thm: The holonomy functor

$$\text{hol}_p : (\mathcal{F}(E), \mathcal{G}(E)) \rightarrow (\text{Hom}(\pi, GL(n, \mathbb{C})), GL(n, \mathbb{C}))$$

is an equivalence of groupoids.

Pf: • surjective on isomorphism classes.

Given a rep $p \in \text{Hom}(\pi, GL(n, \mathbb{C}))$, we construct a flat vector bundle $\mathbb{C}_p \rightarrow \Sigma$ as follows:

the grp π acts on the total space $\sum \times \mathbb{C}^n$ by

$$\gamma \cdot (\tilde{\Sigma}, x) := (\gamma \cdot \tilde{\Sigma}, \underbrace{p(\gamma)x}_{\substack{\uparrow \\ \text{deck transformation}}}) \quad \forall \gamma \in \pi.$$

The quotient $(\sum \times \mathbb{C}^n) / \pi$ is the total space of a smooth vector bundle $\mathbb{C}^n \xrightarrow{P} \Sigma$, which carries a natural flat connection D as the descending of $D_0 = d$ on $\sum \times \mathbb{C}^n$.

$[(\tilde{\Sigma}, \nu)]$ is parallel to D .

So this D gives holonomy P up to conjugation.
 - Full and faithful (need to check)

§4. Rank 1 case for equivalence between Betti and de Rham moduli spaces.

Let E be a trivial complex line bundle over Σ .

A trivialization τ is a global frame of E .

- The gauge transformation $\xi \in G(E)$ is determined by a smooth map $g: \Sigma \rightarrow \mathbb{C}^*$ via $\xi(\tau) = g \cdot \tau$.

$$G(E) \cong \text{Map}(\Sigma, \mathbb{C}^*).$$

$$\text{The subgroup } G_u(E) \cong \text{Map}(\Sigma, U(1))$$

Let $\text{Map}(\Sigma, \mathbb{C}^*)^\circ$ denote the component containing the constant map.

$$G(E)/G(E)^{\circ} = \pi_0(G(E))$$

Note that $\text{Map}(\Sigma, \mathbb{C}^*)^{\circ} \cong \mathcal{A}^0(\Sigma)$

$$g \mapsto \log g.$$

$$\begin{array}{ccc} \log g & \downarrow \exp & \\ \Sigma & \xrightarrow{g} & \mathbb{C}^* \end{array} \quad \text{iff } g_*: \pi_1 \Sigma \rightarrow \pi_1(\mathbb{C}^*) \text{ is trivial}$$

So $\mathcal{F}(E)/G(E) = \frac{(\mathcal{F}(E)/G(E)^{\circ})}{\pi_0(G(E))}.$

- On E , there is a unique connection D_0 s.t. $D_0 \tau = 0$.

Any connection D is of the form

$$D = D_0 + \eta, \quad \eta \in \mathcal{A}^1(\Sigma).$$

D is flat $\Leftrightarrow d\eta = 0$.

$$\mathcal{F}^*(D_0 + \eta) = D_0 + \eta + g^+ dg. \quad (\mathcal{F} \leftrightarrow g \in \text{Map}(\Sigma, \mathbb{C}^*))$$

If $g \in \text{Map}(\Sigma, \mathbb{C}^*)^{\circ}$, $g^+ dg = d \log g$.

So $\mathcal{F}(E)/G(E)^{\circ} \cong Z^1(\Sigma)/B^1(\Sigma) = H^1(\Sigma).$

The Betti moduli space is $\text{Hom}(\pi_1, \mathbb{C}^*) \cong \text{Hom}(\pi_1, S^1) \times \text{Hom}(\pi_1, \mathbb{R}^+)$

The de Rham moduli space:

• $\mathcal{F}(E) = \mathcal{F}_u(E) \times \mathcal{A}^1(\Sigma, \mathbb{R})$

$$\begin{array}{ccc} D_0 + \eta & D_0 + i \text{Im} \eta & \text{Re} \eta \end{array}$$

- $G(E) \cong \text{Map}(\Sigma, \mathbb{C}^*) = \text{Map}(\Sigma, S^1) \times \text{Map}(\Sigma, \mathbb{R}^+)$
 $g \mapsto \begin{matrix} \text{"} \\ G_u(E) \end{matrix} (g_u, g_r)$

$$\begin{aligned} \mathcal{S}^*(D_0 + \eta) &= D_0 + \eta + g^T dg \\ &= (D_0 + i \text{Im} \eta) + \underbrace{(g_u^T \bar{d} g_u)}_{\text{"}} + \underbrace{(g_r^T \bar{d} g_r + \text{Re} \eta)}_{\substack{\text{"} \\ d \log g_r}} \end{aligned}$$

- $\mathcal{F}(E)/G(E) \cong \mathcal{F}_u(E)/G_u(E) \times H^1(\Sigma, \mathbb{R})$
 $\cong \frac{(\mathcal{F}_u(E)/G_u(E))}{\frac{H^1(\Sigma, i\mathbb{R})}{\text{To}(G_u(E))}} \times H^1(\Sigma, \mathbb{R})$

(will see $\cong \frac{H^1(\Sigma, i\mathbb{R})}{H^1(\Sigma, \mathbb{Z})} \times H^1(\Sigma, \mathbb{R})$.)

The equivalence between the moduli spaces is given by

$$\begin{aligned} \text{hol}_p: \mathcal{F}(E) &\rightarrow \text{Hom}(\pi, \mathbb{C}^*) \\ D_0 + \eta &\mapsto (\gamma \mapsto \exp(\int_\gamma \eta)) \end{aligned}$$

Restrict to $\text{hol}_p: \mathcal{F}_u(E) \rightarrow \text{Hom}(\pi, U(1))$

Claim: $\ker(\text{hol}_p) = G(E)$.

Assuming the claim: hol_p descends to a map

$$p: \frac{(\mathcal{F}_u(E)/G_u(E))}{\frac{H^1(\Sigma, i\mathbb{R})}{\text{To}(G_u(E))}} \rightarrow \begin{matrix} \text{Hom}(\pi, U(1)) \\ \text{"} \\ \text{Hom}(\pi, \mathbb{Z}) \\ \text{"} \\ \text{Hom}(\pi, \mathbb{Z}) \end{matrix}$$

$$\Rightarrow \pi_0(G_{\text{un}}(E)) \cong H^1(\Sigma, \mathbb{Z}).$$

$$\cong \mathbb{Z} \langle \omega_1, \dots, \omega_g \rangle$$

where $\omega_1, \dots, \omega_g$ has period $\in 2\pi i \mathbb{Z}$
and form a basis.

Pf of Claim:

" \subset " If $\eta \in H^1(\Sigma, i\mathbb{R})$ s.t. $\exp(\int_{\gamma} \eta) = 1 \quad \forall \gamma \in \pi_1$.

define $g(p) = \exp \int_{x_0}^p \eta$ well-defined.

as a fn $g: \Sigma \rightarrow \mathbb{C}^*$.

and $\eta = g^{-1}dg$.

" \supset " Given $g^{-1}dg$ for $g: \Sigma \rightarrow \mathbb{C}^*$,

one can lift g to $\tilde{g}: \tilde{\Sigma} \rightarrow \mathbb{C}^*$.

then $\tilde{g}^{-1}d\tilde{g}$ is exact on $\tilde{\Sigma}$.

$$\cong \underline{d \log \tilde{g}}.$$

Take a lift $\tilde{\gamma}$ of γ to $\tilde{\Sigma}$,

$$\text{then } \int_{\gamma} g^{-1}dg = \int_{\tilde{\gamma}(0)}^{\tilde{\gamma}(1)} \tilde{g}^{-1}d\tilde{g} = \int_{\tilde{\gamma}(0)}^{\tilde{\gamma}(1)} d \log \tilde{g}$$

$$= \log \tilde{g}(\tilde{\gamma}(1)) - \log \tilde{g}(\tilde{\gamma}(0))$$

$$\in 2\pi i \mathbb{Z}.$$

$$\Rightarrow \exp \int_{\gamma} g^{-1}dg = 1. \quad \square$$

Lecture 2

§5. The Dolbeault groupoid

Let X be a Riemann surface diffeo to Σ .

$$\mathcal{A}'(X) = \mathcal{A}^{1,0}(X) \oplus \mathcal{A}^{0,1}(X)$$

$\quad \quad \quad dz \quad \quad \quad d\bar{z}$

Hodge $*$ -operator on $\mathcal{A}'(X)$

$$*dz = -i dz$$

$$*d\bar{z} = i d\bar{z}.$$

Defn. Given a complex vector bundle E over X ,
a holomorphic structure on E is a diff operator

$$\bar{\partial}_E : \mathcal{A}^{p,q}(X, E) \rightarrow \mathcal{A}^{p,q+1}(X, E)$$

which satisfy

$$\bar{\partial}_E(f \cdot s) = \bar{\partial}f \wedge s + f \cdot \bar{\partial}_E s.$$

$$\forall f \in \mathcal{A}^0(X), s \in \mathcal{A}^{p,q}(X, E)$$

Rmk: If we are dealing w/ higher dim base
mfld, we add the integrability

- The $G(E)$ -action on the space of Higgs bundles is as follows:

$$(\bar{\partial}_E, \Phi) \xrightarrow{\xi} \left(\xi^* \bar{\partial}_E, \xi^* \Phi \right)$$

$$\begin{array}{ccc} & \parallel & \parallel \\ & \xi^* \bar{\partial}_E(\xi \cdot) & \xi^* \Phi \xi \end{array}$$

Locally, w.r.t a frame e ,

$$\bar{\partial}_E e = e \bar{\partial} + \Phi \quad (\bar{\partial}_E = \bar{\partial} + \Phi)$$

$$(\xi^* \bar{\partial}_E) e = e (g^* \bar{\partial} g + g^* \Phi g)$$

$$g \leftrightarrow \xi$$

$$(\xi^* \Phi) e = e \cdot g^* \Phi g$$

Defn. The Dolbeault groupoid is

$$\left(\text{Higgs}(E), G(E) \right)$$

↑
space of Higgs bundles over X .

§ 5.1 Understand the Dolbeault groupoid in $rk=1$, ~~deg 0~~ case.

E is trivial because $rk=1$, $deg=0$ condition.

Start with a trivial complex line bundle E .

A Higgs field on the hol^M line bundle

$(E, \bar{\partial}E)$ is just a hol^M 1-form on X .

since $End(E) = E \otimes E^* = \mathcal{O}$.

• Thus $Higgs(E) = Hol(E) \times \underline{H^{1,0}(X)}$.

\uparrow
space of hol^M str's on E

There is a standard base pt in $Hol(E)$.

$\bar{\partial}_0 = \bar{\partial}$ for $X \times \mathbb{C}$.

An arbitrary hol^M str on E is of the form $\bar{\partial}_0 + \Phi$,

where $\tilde{\psi} \in \mathcal{A}^{0,1}(X)$.

- The gauge action $G(E)$ on $\text{Hol}(E)$
 $\bar{\omega} + \tilde{\psi} \mapsto \bar{\omega} + \tilde{\psi} + \underbrace{g^{-1}\bar{\omega}g}$.

Again, $G(E)/G(E)^{\circ} = \pi_0(G(E))$

$$\text{Hol}(E)/G(E) = \left(\text{Hol}(E)/G(E)^{\circ} \right) / \pi_0(G(E))$$

$$G(E) = \text{Map}(X, \mathbb{C}^*)$$

$$G(E)^{\circ} = \text{Map}(X, \mathbb{C}^*)^{\circ} \quad (\text{containing constant maps})$$

$$g^E = g = \exp(f).$$

So $\text{Hol}(E)/G(E)^{\circ} \cong \mathcal{A}^{0,1}(X) / \text{exact (glt-form)} \bar{\omega} \mathcal{A}^0(X)$.

By the Hodge decomposition

$$H^0(E) / G(E)^0 \cong H^{0,1}(X).$$

So the Dolbeault moduli space

$$\begin{aligned} \text{Higgs}(E) / G(E) &\cong \frac{H^0(E)}{G(E)} \times H^{1,0}(X) \\ &\cong \frac{H^{0,1}(X)}{\underbrace{\pi_0(G(E))}_0} \times H^{1,0}(X). \end{aligned}$$

Claim: $\pi_0(G(E))$'s image form a lattice of $\text{rk } 2g$ in $H^{0,1}(X)$.

(From last time, $\pi_0(G(E))$'s image in $H^1(X)$ is a lattice of $\text{rk } 2g$.)

$$\cong \text{Jac}(X) \times H^{1,0}(X).$$

\uparrow
 a complex torus of dim g .

- Identify $\text{Higgs}(E)/\mathcal{G}(E)$ with $T^*\text{Jac}(X)$.

Consider the Hermitian form on $\mathcal{A}^1(X)$

$$\text{by } \langle \alpha, \beta \rangle := \int_X \alpha \wedge * \bar{\beta}.$$

$\left. \begin{array}{l} \text{(pos. def on } \mathcal{A}^{0,1}(X) \\ \text{(neg. def on } \mathcal{A}^{1,0}(X). \end{array} \right\} d\bar{z} \wedge -i dz.$

Its restriction to $V = H^{0,1}(X)$

defines an isomorphism $\bar{V} \rightarrow V^*$.

of complex v.s.

$$\text{Higgs}(E)/\mathcal{G}(E) = \underbrace{H^{0,1}(X)}_{\cong \text{Jac}(X)} \times \underbrace{H^{1,0}(X)}_{\cong \bar{V}}$$

The tangent space of $\text{Jac}(X)$ at any pt identifies with V .

$$\text{Thus } V/L \times \bar{V} \cong V/L \times V^*$$

$$\Rightarrow \text{Higgs}(E)/G(E) \cong T^*\text{Jac}(X).$$

§ 6. Equivalence between the de Rham and Dolbeault groupoids for $\text{rk } 1, \text{deg } 0$ case.

§ 6.1. Introduce Hermitian metrics.

Defn. A Herm metric H on E is a smooth family of pos. def Herm forms $\langle \cdot, \cdot \rangle_H: E_x \times E_x \rightarrow \mathbb{C}$.

Denote by $\text{Her}(E)$ the space of Hermitian metrics on E .

In terms of a basis (frame) e ,

$$H(\underline{e \cdot \xi}, \underline{e \cdot \eta}) = \xi^t h \eta \quad \text{Hermitian matrix.}$$

where $h = H(e, e)$

• The action of $G(E)$ on $\text{Her}(E)$,
locally, $g \cdot h = (g^{-1})^t h g^{-1}$.

• If E is a flat vector bundle on X with holonomy $\phi: \pi \rightarrow G(r, \mathbb{C})$, then a Hermitian metric $H \in \text{Her}(E)$ corresponds to

a ϕ -equivariant map

$$h: \overset{\vee}{X} \longrightarrow \text{Her}(\mathbb{C}^r)$$

(i.e. $h(\sigma \cdot x) = \phi(\sigma) h(x)$.)

Idea: Parallel transport H along paths based at x_0 to E_{x_0} w.r.t this flat connection.

i.e. $h([\sigma])^{(u,u)} = H(\underline{s(\sigma(t))}, s(\sigma(t)))$

where $s(\sigma(t))$ is a parallel section along σ starting from u .

- Induced Hermitian pairing over $\mathcal{A}^*(X, E)$
 $\mathcal{A}^k(X, E) \times \mathcal{A}^l(X, E) \rightarrow \mathcal{A}^{k+l}(X).$

Defn. A connection D is unitary w.r.t H if

$$d \langle S_1, S_2 \rangle_H = \langle DS_1, S_2 \rangle_H + \langle S_1, DS_2 \rangle_H.$$

Prop. Given $(E, \bar{\partial}_E)$ with H ,
 $\exists!$ a connection D st

$$(1) D^{0,1} = \bar{\partial}_E.$$

(2) D is unitary w.r.t H .

D is called Chern connection.

Prop. Given a connection D and H ,
 $\exists!$ a decomposition

$$D = D_H + \tilde{\Psi} \lrcorner H \quad \left\{ \begin{array}{l} \text{self-adjoint} \\ \text{w.r.t } H. \end{array} \right.$$

st unitary connection
w.r.t H .

$$\left(H(\tilde{\Psi} \lrcorner S, t) := \frac{1}{2} \{ H(Ds, t) + H(S, D(t)) - d(H(s, t)) \} \right)$$

§ 6.2. Restrict to the case

E is a complex trivial line bundle.
with a frame τ .
trivialization.

Let H_0 be $H_0(\tau, \tau) = 1$.

$$G(E) = \text{Map}(X, \mathbb{C}^*) \ni g.$$

• $G(E)$ acts on $\text{Her}(E)$ as
 $h \mapsto |g|^2 h$.

$$\left(\langle u, v \rangle_{\mathbb{C} \cdot H} := \langle g^{-1}u, g^{-1}v \rangle_H \right)$$

$$h = H(\tau, \tau) : X \rightarrow \mathbb{R}^+$$

• $G(E)$ acts on $\text{Her}(E)$ transitively.

Want $g \cdot h_1 = h_2$,

$$\text{need } g(z) = \sqrt{\frac{h_1(z)}{h_2(z)}}.$$

- $D = D_0 + \eta$ is unitary w.r.t H
 iff
 $d(H(\tau, \tau)) = H(D\tau, \tau) + H(\tau, D\tau)$
 $\Rightarrow dh = h\eta + h\bar{\eta}$
 $\Rightarrow h^{-1}dh = \eta + \bar{\eta} = 2\text{Re}(\eta).$

- Ψ is self-adjoint w.r.t H
 iff $\Psi = \bar{\Psi}$ i.e. Ψ is real.
 $(H(\Psi\tau, \tau) = H(\tau, \Psi\tau))$

- w.r.t H , $D = D_0 + \eta$ is uniquely decomposed $D = D_H + \Psi_H$,
 $\begin{cases} D_H = D_0 + i\text{Im}\eta + \frac{1}{2}h^{-1}dh. \\ \Psi_H = \text{Re}\eta - \frac{1}{2}h^{-1}dh. \end{cases}$
-

§6.2(a). Start from a flat connection

Goal: To find a decomposition and further obtain a Higgs bundle.

Idea: Find the "best" H and use H to decompose.

Defn. A Hermitian metric h is harmonic w.r.t $D \in \mathcal{F}(E)$

if the associated equivariant map $h: X \rightarrow \text{Herm}(\mathbb{C}) = \mathbb{R}^t$.

(*) corresponding to h

is a multiplicatively harmonic fun.

(defined as, its logarithm is a harmonic fun.)

Prop. Condition $(*)$ holds

iff ζ_H is a harmonic 1-form.

(iff $(\zeta_H)^{1,0}$ is holomorphic.)

Pf: For a path $\gamma: [0,1] \rightarrow X$ with $\gamma(0) = x_0$.

Let s be the flat section from τ along γ , then

$$(s(x_0) = \tau) \quad s(\gamma(1)) = \exp\left(-\int_{\gamma} \eta\right) \cdot \tau_0$$

$$\text{for } D = D_0 + \eta.$$

Use the defn of \tilde{h} ,

$$\tilde{h}([\gamma]) := H(s(\gamma(1)), s(\gamma(0)))$$

$$\begin{array}{c} \uparrow \\ X \end{array} = H(\exp(-\int_{\gamma} \eta) \cdot \tau_0, \exp(-\int_{\gamma} \eta) \cdot \tau_0)$$

$$\begin{aligned}
&= \exp\left(-\int_{\gamma} (\eta + \bar{\eta})\right) \cdot \frac{H(\tau_0, \tau_0)}{h} \\
&= \exp\left(-2\int_{\gamma} \operatorname{Re}(\eta) + h^{-1}dh\right) \\
&= \exp\left(-2\int_{\gamma} \left(\operatorname{Re}(\eta) - \frac{1}{2}h^{-1}dh\right)\right) \\
&= \exp\left(-2\int_{\gamma} \Psi_H\right). \quad \square
\end{aligned}$$

Prop. For $D \in \mathcal{F}(E)$, $\exists!$ a harmonic metric H (up to constant scalar).

$$\begin{aligned}
\text{Pf: } D &= D_H + \Psi_H \\
&= (D_0 + i\operatorname{Im}\eta) + \frac{1}{2}h^{-1}dh \\
&\quad + \frac{\operatorname{Re}\eta - \frac{1}{2}h^{-1}dh}{}
\end{aligned}$$

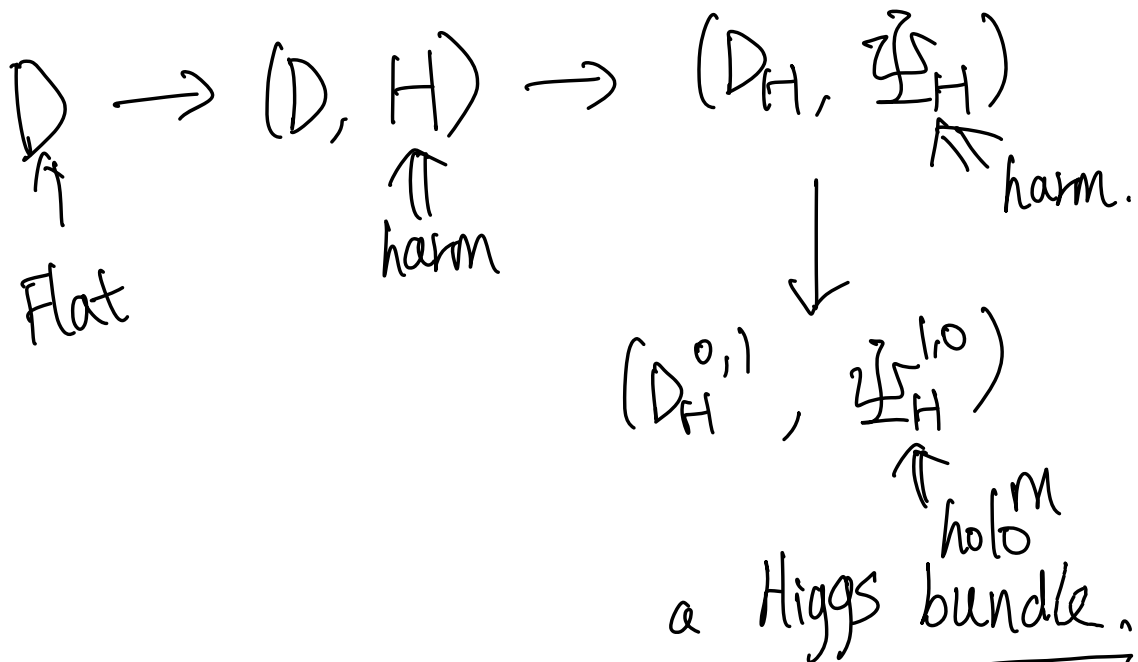
One wants to find h s.t

$\operatorname{Re} \eta - \frac{1}{2} h^{-1} dh$ is harmonic.

By Hodge decomposition,

$$\operatorname{Re} \eta = \text{harmonic 1-form} + \underline{dS}$$

Let $h = e^{2S}$.



§ 6.2(b). Start from a Higgs bundle

One wants to find a way to sum up and obtain a flat connection.

Defn. A Hermitian metric is harmonic w.r.t $\bar{\partial}_E$ iff the Chern connection D^H is flat. iff $D^H + \phi + \bar{\phi}$ is flat.

Prop. For each $\bar{\partial}_E \in \text{Hol}(E)$, $\exists!$ $h \in \text{Her}(E)$ up to constant s.t h is harmonic w.r.t $\bar{\partial}_E$.

Pf: Write $\bar{\partial}_E = \bar{\partial}_0 + \bar{\psi}$
 $= \bar{\partial}_0 + (\bar{\psi}_0 + \bar{\alpha}s)$
 $\quad \quad \quad \uparrow$
 $\quad \quad \quad H^{0,1}(X)$

Note that H_0 is harmonic w.r.t $\bar{\partial}_0 + \bar{\psi}_0$,
 $(D^{H_0} = D_0 + \bar{\psi}_0 - \bar{\psi}_0)$

Then $g \cdot H_0$ is harmonic to $g \cdot (\bar{\partial}_0 + \bar{\psi}_0)$.

Let $g = e^S$,

$$g^{-1} \cdot (\bar{\alpha}_0 + \bar{\Phi}_0) = \bar{\alpha}_0 + \bar{\Phi}_0 + \bar{\alpha}S = \bar{\alpha}_E$$

So $g^{-1} \cdot H_0$ is the desired metric.

$$\stackrel{''}{e^{2S}} H_0$$

□

§ 6.2 (c) Combine

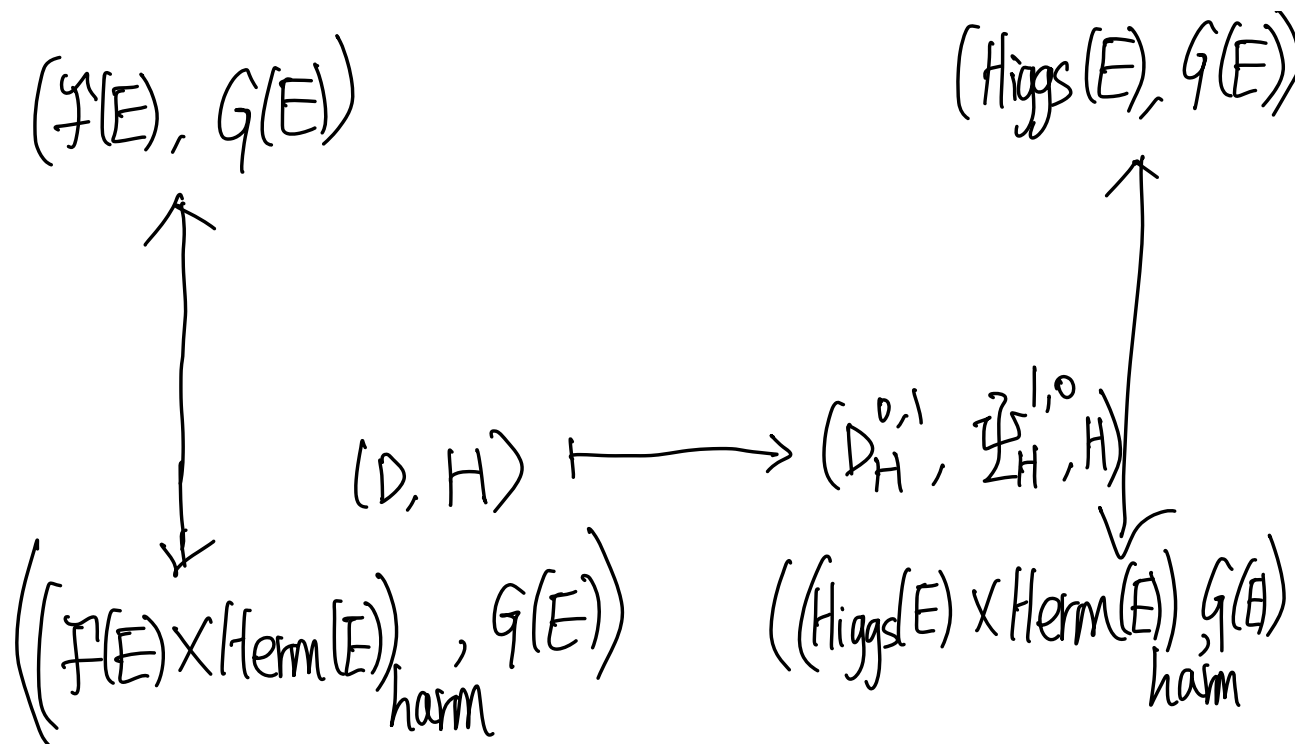
Denote by $(\mathcal{F}(E) \times \text{Her}(E))_{\text{harm}}$ the subset

of (D, H) s.t. H is harm w.r.t. D

Denote by $(\text{Higgs}(E) \times \text{Her}(E))_{\text{harm}}$ the subset

of $(\bar{\alpha}_E, \phi, H)$ s.t. H is harm w.r.t. $\bar{\alpha}_E$.

We have a diagram:




$$(D^H + \phi + \bar{\phi}, H) \longleftarrow (\bar{\partial}_E, \phi, H)$$

Thm. The induced functor

$$(\mathcal{F}(E), G(E)) \longrightarrow (\text{Higgs}(E), G(E))$$

is an equivalence of groupoids.

Pf: The rest need to check. 

§7. Complex structures on moduli spaces. (again rk 1 cases.)

- Betti moduli space

$$M_{\text{Betti}} = \text{Hom}(\pi, \mathbb{C}^*) \cong (\mathbb{C}^*)^{2g}$$

$$T_p \text{Hom}(\pi, \mathbb{C}^*) \cong \text{Hom}(\pi, \mathbb{C}) = \mathbb{C}^{2g}.$$

$$J_1: T_p M_{\text{Betti}} \rightarrow T_p M_{\text{Betti}}$$

$$X \mapsto iX.$$

- De Rham moduli space.

$$M_{\text{de Rham}} \cong \frac{H^1(X, i\mathbb{R}) \times H^1(X, \mathbb{R})}{H^1(X, \mathbb{Z})}$$

$$\cong H^1(X) / H^1(X, \mathbb{Z})$$

$$T_\eta M_{\text{de Rham}} \cong H^1(X)$$

$$J_2: T_\eta M_{\text{de Rham}} \rightarrow T_\eta M_{\text{de Rham}}$$

$$X \mapsto iX.$$

In fact, the equivalence between

$$\begin{aligned} \mathcal{M}_{\text{de Rham}} &\longrightarrow \mathcal{M}_{\text{Betti}} \\ D_0 + \eta &\longmapsto (\gamma \mapsto \exp \int_{\gamma} \eta) \end{aligned}$$

The tangent map is at η ($X \mapsto (\gamma \mapsto \exp \int_{\gamma} \eta, \int_{\gamma} X)$)

is a biholomorphism w.r.t $\mathcal{J}_1, \mathcal{J}_2$.

So we can say $\mathcal{J}_1 = \mathcal{J}_2$, denoted by \mathcal{J} .

- Dolbeault moduli space : $\mathcal{M}_{\text{Dol}}, \mathcal{M}_{\text{Higgs}}$

$$\mathcal{M}_{\text{Dol}} = \frac{H^{0,1}(X) \times H^{1,0}(X)}{\text{Jac}(X)}$$

$$T_0 \mathcal{M}_{\text{Dol}} \cong H^{0,1}(X) \times H^{1,0}(X) \\ (\Psi, \Phi)$$

$$I: T_0 \text{Mod} \rightarrow \overline{T_0 \text{Mod}}$$
$$(\Phi, \Psi) \mapsto (i\Phi, i\Psi).$$

I is different from J .

Lecture 3

X — a compact R.S
of $g \geq 2$ if not specified.

$\mathcal{E} — (E, \bar{\partial}_E)$ holo^m v.b.

Goal today:

- moduli space of polystable
vector bundles
/ Higgs bundles
- relate the stability with
soln to Hitchin eqn.

§1. Preparation: extensions of holo^m vector bundles.

If $\mathcal{F} \subset \mathcal{E}$ is a holo^m subbundle, then $Q = \mathcal{E}/\mathcal{F}$ has an induced holo^m str, $\bar{\partial}Q$.

Moreover, we can choose a complement subbundle of \mathcal{F} inside \mathcal{E} to represent Q .
 (e.g. choose $Q = \mathcal{F}^\perp$)

Write $\mathcal{E} = \mathcal{F} \oplus Q$ C^∞ direct sum

$$\text{Then } \bar{\partial}\mathcal{E} = \begin{pmatrix} \bar{\partial}\mathcal{F} & \beta \\ 0 & \bar{\partial}Q \end{pmatrix},$$

where $\beta \in A^{0,1}(X, \text{Hom}(Q, \mathcal{F}'))$,

called the second fundamental form.

- If β is in $\bar{\omega}(A^0(X, \underline{\text{Hom}}(Q, F)))$,
then by a gauge transformation,
 $E = F \oplus Q'$ hol^m splitting.

- Isomorphism classes of E of
the form $0 \hookrightarrow F \rightarrow E \rightarrow Q \rightarrow 0$ (*)
is in bijection with
 $\mathbb{P}(H_{\bar{\omega}}^{0,1}(X, \underline{\text{Hom}}(Q, F)))$

- Call the extension sequence (*) split
if $[\beta] = 0$.

iff \exists an injection $Q \hookrightarrow E$
lifting the proj $E \rightarrow Q$.

- For such splitting $E = F \oplus Q$,
 $\bar{\omega}_E = \begin{pmatrix} \bar{\omega}_F & \beta \\ 0 & \bar{\omega}_Q \end{pmatrix}$

If $H = \begin{pmatrix} H_1 & \\ & H_2 \end{pmatrix}$ w.r.t $E = F \oplus Q$,
 then the Chern connection ∇^H determined
 by $\bar{\partial}E, H$ is

$$\nabla^H = \begin{pmatrix} \nabla^F & \beta \\ -\beta^{*h} & \nabla^Q \end{pmatrix},$$

where ∇^F, ∇^Q are the Chern connection
 β^{*h} is the adjoint of β .
 $(1,0)$ -form $(0,1)$ -form

The curvature of ∇^H is

$$F(\nabla^H) = \begin{pmatrix} F(\nabla^F) - \beta \wedge \beta^{*h} & \partial\beta \\ -\bar{\partial}\beta^{*h} & F(\nabla^Q) - \beta^{*h} \wedge \beta \end{pmatrix}$$

§2. Moduli space of hol^m vector bundles

key: To introduce stability on hol^m v.b.

Two motivations for stability:

① Original motivation due to Mumford is to provide the set of gauge equivalence classes of hol^m v.b with a "good" topology.

Mainly, the unstable ones cause "non-Hausdorff" problem.

② Turns out stability is an iff condition for a hol^m v.b admitting a soln to the Hermitian-Einstein eqn.

Let $\mu(E)$ denote the slope of E ,
ie. $\text{deg}(E) / \text{rank}(E)$.

Defn (Mumford)

- A holomorphic v.b. E is called stable **semistable** if $\mu(F) < \mu(E)$ for any proper holomorphic subbundle F of E .
- A holomorphic v.b. E is called polystable if it is a direct sum of stable holomorphic subbundles of the same slope.

Remark: (1) Stability is preserved under gauge transformations.
(2) Stability is an open condition.

For an unstable vector bundle,

Prop. Given an arbitrary hol^m v.b.,

$\exists!$ a Harder-Narasimhan filtration of E ,

$$0 = E_0 \subset E_1 \subset E_2 \subset \dots \subset E_p = E$$

st. • E_i/E_{i-1} are semistable.

• $\mu(E_i/E_{i-1})$ is strictly decreasing.

(Idea: Take E_1 to be the maximal destabilizing subbundle of E .)

For a semistable v.b.,

Prop. Given a semistable v.b.,

\exists a Jordan-Hölder filtration of E ,

$$0 = E_0 \subset E_1 \subset E_2 \subset \dots \subset E_p = E$$

st. the quotients E_i/E_{i-1} are stable

Obviously, $\mu(E_i/E_{i-1})$ are the same.

Denote by $\text{Gr}(E) = \bigoplus_{i=1}^n E_i/E_{i-1}$,
 (graded v.b of E)
 it is polystable.

Defn. Two semistable v.b are S -equiv
 if their graded v.b's are gauge equiv.

Rmk: $\{\text{stable}\} \subset \{\text{polystable}\} \subset \{\text{semistable}\}$.

Denote $M^S(r, d) := \left\{ \begin{array}{l} \text{stable holo v.b of} \\ \text{rk } r, \text{ deg } d \end{array} \right\} / G$

$M(r, d) := \left\{ \begin{array}{l} \text{polystable holo v.b of} \\ \text{rk } r, \text{ deg } d \end{array} \right\} / G$

$\cong \left\{ \begin{array}{l} \text{semistable holo v.b of} \\ \text{rk } r, \text{ deg } d \end{array} \right\} / S\text{-equiv}$

Note that when $(r, d) = 1$, $M^S(r, d) = M(r, d)$.

And $N^S(r, d)$ is a smooth cpt complex mfd.

Ex 1. On \mathbb{P}^1 , by Grothendieck's thm,
any holom v.b is of the form
$$E = \bigoplus_{i=1}^k \mathcal{O}(n_i).$$

$$\text{Then } \mu(E) = \frac{\sum_{i=1}^k n_i}{k}.$$

E is unstable unless all n_i 's are equal.
polystable if all n_i 's are equal.
stable only if $k=1$.

Ex 2. Consider the extension sequence

$$0 \rightarrow \mathcal{O} \rightarrow \mathcal{E} \rightarrow \mathcal{O}(p) \rightarrow 0.$$

The isomorphism classes of \mathcal{E} are
parametrized by $\mathbb{P}(H_{\mathbb{C}}^{0,1}(X, \mathcal{O}(-p)))$

$$= \mathbb{P}(H^0(X, K(P))^*),$$

has $\dim_{\mathbb{C}} = g$.

Claim = Any non-split extension of this type is stable.

Pf: If $L \hookrightarrow E$ is a destabilizing line subbundle,

$$\left(\deg L > \frac{\deg(E)}{\operatorname{rk}(E)} = \frac{1}{2}. \right)$$

then $\deg L \geq 1$.

Can compose $L \hookrightarrow E \rightarrow \mathcal{O}(P)$,

obtain a holom map $L \rightarrow \mathcal{O}(P)$,

which is either 0 or an isomorphism.

(i) If $L \rightarrow \mathcal{O}(P)$ is 0,

then $L = \ker(L \rightarrow \mathcal{O}(P))$

$$C \subset \ker(\mathcal{E} \rightarrow \mathcal{O}(p))$$

$$\mathcal{O} \cdot \quad \swarrow \searrow$$

(ii) If $\mathcal{L} \rightarrow \mathcal{O}(p)$ is an isom,

then $\mathcal{O}(p) \rightarrow \mathcal{L} \hookrightarrow \mathcal{E}$ is a nontrivial map

lifting the projection $\mathcal{E} \rightarrow \mathcal{O}(p)$.

So the extension splits. $\swarrow \searrow \square$

Note that $\mathcal{E} = \mathcal{O} \oplus \mathcal{O}(p)$ is unstable.

Ex 3. Claim: $0 \rightarrow \mathcal{O}_1 \rightarrow \mathcal{E} \rightarrow \mathcal{O}_2 \rightarrow 0$
with non-split extension
is strictly semistable.

Pf: • Not stable, since \mathcal{O}_1 has the same slope 0 as \mathcal{E} .

• semistable:

For any $L \subset E$ a holomorphic line subbundle,
 The induced map $L \rightarrow \mathcal{O}_2$ is either zero
 or nontrivial.

(i) If $L \rightarrow \mathcal{O}_2$ is 0, then $L \hookrightarrow \mathcal{O}_1$.
 $\Rightarrow \deg L \leq 0$.

(ii) If $L \rightarrow \mathcal{O}_2$ is nontrivial, $\Rightarrow \deg L \leq 0$. \square

Note that $\text{Gr}(E) = \mathcal{O}_1 \oplus \mathcal{O}_2$.

Claim: $\mathcal{O}_1 \oplus \mathcal{O}_2$ is contained in
 holomorphic direct sum
 the closure of the gauge orbit of
 E with non-split extension.

Pf: Take the 1-parameter subgroup
 $g_t = \begin{pmatrix} t & 0 \\ 0 & t^{-1} \end{pmatrix} (t > 0)$.

Then $g_t^* \bar{\partial} E = g_t^* \left(\bar{\partial} + \begin{pmatrix} 0 & \beta \\ 0 & 0 \end{pmatrix} \right)$

$$\begin{aligned}
&= \bar{\partial} + g_t^{-1} \begin{pmatrix} 0 & \beta \\ 0 & 0 \end{pmatrix} g_t + \underbrace{g_t^{-1} \bar{\partial} g_t}_{\text{0}} \\
&\quad \begin{matrix} \text{"} \\ (t^{-1} \quad t) \end{matrix} \quad \begin{matrix} \text{"} \\ (t \quad t^{-1}) \end{matrix} \\
&= \bar{\partial} + \begin{pmatrix} 0 & t^2 \beta \\ 0 & 0 \end{pmatrix} \\
&= \bar{\partial} \quad \text{as } t \rightarrow \infty. \quad \square
\end{aligned}$$

Let ω_X be a Kähler form on X
to be normalized s.t. $\int_X \omega_X = 1$.

Thm. (Narasimhan - Seshadri)

A hol^m v.b E carries a Herm
metric h satisfying the Hermitian-Einstein equ

$$F(\nabla h) = -2\pi i \cdot \mu(E) \cdot \text{id}_E \cdot \omega_X$$

\uparrow
 $\mathcal{A}^2(\text{End}(E))$

iff E is polystable.

Moreover, the soln h is unique up to multiplication by a positive constant if E is stable.

Remark: From Chern-Weil theory,

$$\frac{i}{2\pi} \int_X \text{Tr}(F(\sigma)) = \text{deg}(E)$$

Use the eqn, for any connection on E .

$$\frac{i}{2\pi} \int_X \text{Tr}(F(\nabla^h)) = \frac{i}{2\pi} \int_X -2\pi i \cdot \mu(E) \cdot \text{Tr}(\text{id}_E) \cdot \omega_X$$

$$\parallel$$
$$\mu(E) \cdot \text{rk}(E) \cdot \int_X \omega_X$$

$$\parallel$$
$$\text{deg}(E).$$

Rmk: The first proof to N-S was algebraic and relates stable v.b with unitary reps of $\pi = \pi_1(X)$. Donaldson presented an analytic proof. N-S thm holds for cpt Kähler mflds of arbitrary dim, which is the Donaldson-Uhlenbeck-Yau thm.

§3. Moduli space of Higgs bundles

Repeat all the defs to Higgs bundles.

Defn. A Higgs bundle (E, ϕ) is stable semistable

if $\mu(F) \leq \mu(E)$ for any proper ϕ -inv holom subbundle F of E .

polystable if $(E, \phi) = \bigoplus_i (E_i, \phi_i)$
 E_i stable w/ the same slope.

- H-N filtration for Higgs bundles
- J-H filtration for semistable Higgs bundle
graded polystable Higgs bundles

$$\mathcal{M}^{\text{Higgs}, S}(r, d) = \left\{ \begin{array}{l} \text{stable Higgs bundles} \\ \text{of rk } r, \text{ deg } d \end{array} \right\} / G$$

$$\mathcal{M}^{\text{Higgs}}(r, d) = \left\{ \begin{array}{l} \text{polystable Higgs bundles} \\ \text{of rk } r, \text{ deg } d \end{array} \right\} / G$$

$$\cong \left\{ \begin{array}{l} \text{semistable Higgs bundles} \\ \text{of rk } r, \text{ deg } d \end{array} \right\} / S\text{-equiv}$$

- $\{\text{stable}\} \subset \{\text{polystable}\} \subset \{\text{semistable}\}$

- $(r, d) = 1, \mathcal{M}^{\text{Higgs}}(r, d) = \mathcal{M}^{\text{Higgs}, S}(r, d)$

For our interests, we also focus on $SL(n, \mathbb{C})$ -Higgs bundles.

Defn. An $SL(n, \mathbb{C})$ -Higgs bundle is a Higgs bundle (E, ϕ)

s.t. • $\det(E) \cong \mathcal{O}$

• $\text{tr}(\phi) = 0$.

Rmk: When we construct moduli space of $SL(n, \mathbb{C})$ -Higgs bundles, polystable

the gauge transformation lies in $SL(n, \mathbb{C})$.

Rmk: One can even consider G -Higgs bundles for reductive Lie grps G .

e.g. $G = SO(n, \mathbb{C}), SL(n, \mathbb{R}), Sp(2n, \mathbb{R}), \dots$
and associated stability.

Ex. Fix a holomorphic line bundle $K^{\frac{1}{2}}$ of K
($(K^{\frac{1}{2}})^2 = K$.)

Let $\mathcal{E} = K^{\frac{1}{2}} \oplus K^{-\frac{1}{2}}$ holomorphic direct sum.

Claim: The Higgs bundle (\mathcal{E}, ϕ_q)

with $\phi_q = \begin{pmatrix} 0 & q \\ 1 & 0 \end{pmatrix}: \mathcal{E} \rightarrow \mathcal{E} \otimes K$

is stable, where $q \in H^0(X, K^{\otimes 2})$.

Pf: (i) $q=0$ case. $\phi_0 = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$

The only ϕ_0 -inv holomorphic subbundle

is $K^{-\frac{1}{2}}$, which has $\deg 1 - g < 0$
" " $\deg(\mathcal{E})$

So it is stable.

(ii) Use the fact that stability is an open condition the Higgs bundle

$(\mathcal{E}, \Phi_{\varepsilon q} = \begin{pmatrix} 0 & \varepsilon q \\ 1 & 0 \end{pmatrix})$ is stable.

(iii) Use $g = \begin{pmatrix} \varepsilon^{\frac{1}{4}} & \\ & \varepsilon^{-\frac{1}{4}} \end{pmatrix}$

Then (ii) means

$(\mathcal{E}, g^{-1} \Phi_{\varepsilon q} g)$ is again stable

$$\begin{pmatrix} \varepsilon^{-\frac{1}{4}} & \\ & \varepsilon^{\frac{1}{4}} \end{pmatrix} \begin{pmatrix} 0 & \varepsilon q \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \varepsilon^{\frac{1}{4}} & \\ & \varepsilon^{-\frac{1}{4}} \end{pmatrix}$$

$$\varepsilon^{\frac{1}{2}} \begin{pmatrix} 0 & q \\ 1 & 0 \end{pmatrix}$$

Claim: If (E, ϕ) is stable, so is $(E, t\phi)$ $\forall t \in \mathbb{C}^*$

Pf: Obvious. \square

§4. Hitchin-Kobayashi correspondence

Defn. Given a Higgs bundle (E, ϕ) over X ,
call a Herm metric h harmonic

if it solves the Hitchin eqn
(Hitchin's self-duality eqn)

$$F(\nabla^h) + [\phi \wedge \phi^{*h}] = -2\pi i \cdot \mu(E) \cdot \text{id}_E$$

$\cdot \omega_X$

where

- ∇^h is the Chern connection.

- $F(\nabla^h)$ is the curvature.

- ϕ^{*h} is the adjoint of ϕ w.r.t h ,

ie. $h(\phi s, t) = h(s, \phi^{*h} t)$.

Rmk: (i) Locally, if $\underline{\Phi} = \varphi d\bar{z}$ w.r.t some frame of E .

The metric presentation is h locally.

$$\text{Then } \underline{\Phi}^{*h} = \varphi^{*h} d\bar{z}$$

$$(\varphi^t h = h \overline{\varphi^{*h}} \Rightarrow \varphi^{*h} = h^{-1} \overline{\varphi^t h}.)$$

$$\text{Thus } [\underline{\Phi}, \underline{\Phi}^{*h}] = [\varphi, \varphi^{*h}] dz \wedge d\bar{z}$$

$$\text{Or globally, } [\underline{\Phi}, \underline{\Phi}^{*h}] = \phi \wedge \phi^{*h} + \phi^{*h} \wedge \phi.$$

(ii) When $\deg(E) = 0$ ($\mu(E) = 0$),

Claim: the Hitchin eqn is equivalent to

$$D = \nabla^h + \phi + \phi^{*h} \text{ is flat.}$$

Then we obtain a map from

{Higgs bundles which admits harmonic metric}

\longrightarrow {flat connections} \xrightarrow{G}

Pf of Claim: $F(D) = \underbrace{F(\nabla^h) + [\phi, \phi^{*h}]}_{\text{skew-Her}} + \underbrace{\nabla^h(\phi + \phi^{*h})}_{\text{Herm}} = 0$

$$\Leftrightarrow \begin{cases} F(\nabla^h) + [\phi, \phi^{*h}] = 0 \\ \nabla^h(\phi + \phi^{*h}) = 0. \end{cases}$$

$$\Leftrightarrow \begin{cases} \text{Next time!} \end{cases}$$